Spherical Collapse Model with Clustered Dark Energy

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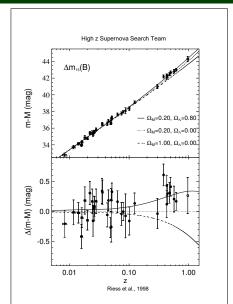
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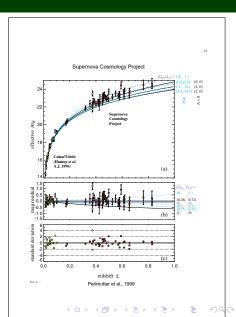
Jun., 12, 2018

Outline

- Accelerating Expansion of the Universe
- Structure Formation
- Virialization
- Complete Journey of Matter Density Perturbation
- Structure Formation with Time-Varying Equation of State
- Conclusions

Type la Supernova Data





According to the cosmological principle and the observation, the universe is homogeneous and isotropic on large scales (> 100 Mpc).

FRW metric:

$$ds^2 = -dt^2 + \underline{a(t)}^2 \left(\frac{dr^2}{1 - \underline{k}r^2} + r^2 d\Omega^2 \right)$$

scale factor curvature of the space

⇒ a spatially homogeneous and isotropic universe expanding as a function of time



Friedmann Equations

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2}$$
$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P)$$

if $P<-rac{1}{3}
ho\Rightarrow\ddot{a}>0$: acceleration (Dark Energy)

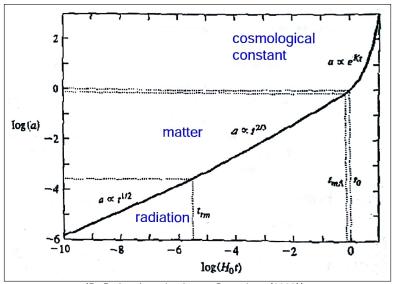
Fluid Equation

energy-momentum conservation:

$$\dot{\rho} + 3\frac{\dot{a}}{a}(\rho + P) = 0$$

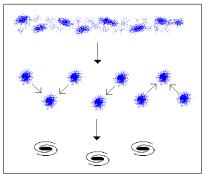
for
$$P = w\rho \Rightarrow \rho(a) = \rho_0 a^{-3(1+w)}$$





(B. Ryden, Introduction to Cosmology (2002))

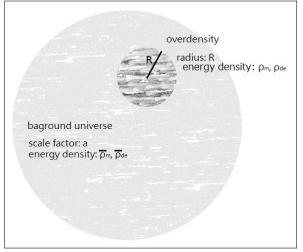
Bottom-Up Structure Formation:
 The small galaxies form and attract each other by gravity and merge to form larger galaxies, and then cluster together to form clusters.

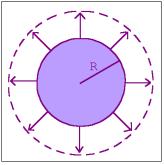


 $(https://www.astro.ufl.edu/\!\sim\!guzman/ast1002/class_notes/Ch15/Ch15b.html)$

Gravity \longleftrightarrow Universe Expansion







- expand to $R_{\rm max}$
- turn around
- collapse to form structures
- ⇒ We use spherical collapse model to explore the nonlinear gravitational collapse in matter.

Background Universe

$$egin{align} rac{\ddot{a}}{a} &= -rac{4\pi G}{3} \left[ar{
ho}_{
m m} + (1+3w) ar{
ho}_{
m de}
ight] \ & \ \dot{ar{
ho}}_{
m m} + 3 \left(rac{\dot{a}}{a}
ight) ar{
ho}_{
m m} = 0 \ & \ \dot{ar{
ho}}_{
m de} + 3 (1+w) \left(rac{\dot{a}}{a}
ight) ar{
ho}_{
m de} = 0 \ & \ \end{aligned}$$

Spherical Overdensity

$$\begin{split} \frac{\ddot{R}}{R} &= -\frac{4\pi G}{3} \left[\rho_{\rm m} + (1+3w) \rho_{\rm de} \right] \\ \dot{\rho}_{\rm m} + 3 \left(\frac{\dot{R}}{R} \right) \rho_{\rm m} = 0 \\ \dot{\rho}_{\rm de} + 3(1+w) \left(\frac{\dot{R}}{R} \right) \rho_{\rm de} = \alpha \Gamma, \\ \text{where } \Gamma = 3(1+w) \left(\frac{\dot{R}}{R} - \frac{\dot{a}}{a} \right) \rho_{\rm de} \\ \text{with } 0 \leq \alpha \leq 1 \end{split}$$

$$\dot{\rho}_{\mathrm{de}} + 3(1+w)\left(\frac{\dot{R}}{R}\right)\rho_{\mathrm{de}} = \alpha\Gamma,$$

where
$$\Gamma=3(1+w)\left(rac{\dot{R}}{R}-rac{\dot{a}}{a}
ight)
ho_{\mathrm{de}}$$
 with $0\leqlpha\leq1$

 $\alpha = 1$ (non-clustering DE)

$$\dot{
ho}_{
m de} + 3(1+w)\left(rac{\dot{a}}{a}
ight)
ho_{
m de} = 0 \Rightarrow
ho_{
m de} = ar{
ho}_{
m de}$$

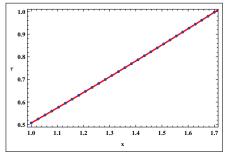
- DE is homogeneous.
- Energy does not conserve within the overdensity.

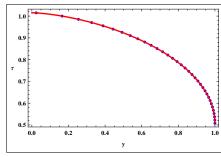
$\alpha = 0$ (clustering DE)

$$\dot{
ho}_{
m de} + 3(1+w)\left(rac{\dot{R}}{R}
ight)
ho_{
m de} = 0 \Rightarrow
ho_{
m de}
eq ar{
ho}_{
m de}$$

- DE is inhomogeneous.
- The spherical overdensity is considered an isolated system satisfying the law of energy conservation.

non-clustering DE (lpha=1), $ho_{
m de}=ar{
ho}_{
m de}$ (Lee and Ng, JCAP 10 (2010) 028)





•
$$w = -\frac{1}{3}$$
, $z_{\rm ta} = -0.8$, $\Omega_{\rm m,0} = 0.3$

$$*d au \equiv extit{H}_{
m ta}\sqrt{\Omega_{
m m}(extit{x}_{
m ta})}dt$$
 , $extit{x} \equiv rac{ extit{a}}{ extit{a}_{
m ta}}$, $extit{y} \equiv rac{ extit{R}}{ extit{R}_{
m ta}}$

clustering DE ($\alpha = 0$), $\rho_{\rm de} \neq \bar{\rho}_{\rm de}$

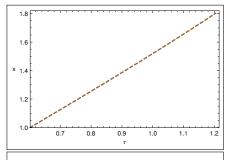
- Difficulties: $\rho_{\rm de}$ within the spherical overdensity is unknown.
- \Rightarrow model DE by various types of scalar fields
 - Creminelli et al., JCAP 03 (2010) 027
 - Nunes and Mota, MNRAS 368 (2006) 751
 - ▶ Mota and Bruck, Astron. Astrophys 421 (2004) 71
- ⇒ Might it oversimplify the property of DE? (e.g. early DE model)
 - ► Huey et al., PRD 59 (1999) 063005
 - Doran and Lilley, MNRAS 330 (2001) 965
 - ▶ Bean et al., PRD 64 (2001) 103508
 - Lee and Ng, PRD 67 (2003) 107302

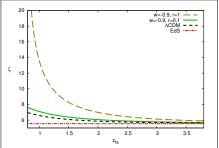


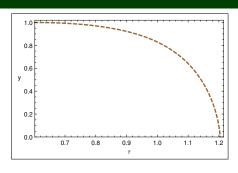
clustering DE ($\alpha=0$), $\rho_{\mathrm{de}} \neq \bar{\rho}_{\mathrm{de}}$

$$\begin{array}{lll} \frac{dx}{d\tau} & = & \sqrt{x^{-1} + \frac{1}{Q_{\mathrm{ta}}} x^{-3w-1}} \\ & \Rightarrow & \tau_{\mathrm{ta}} = \frac{2}{3} F\left[\frac{1}{2}, -\frac{1}{2w}, 1 - \frac{1}{2w}, -Q_{\mathrm{ta}}^{-1}\right] \\ \frac{d^2y}{d\tau^2} & = & -\frac{1}{2} \left[\zeta y^{-2} + (1+3w) \frac{1}{Q_{\mathrm{cta}}} y^{-3(1+w)+1} \right] \\ \zeta: & \rho_{\mathrm{m}} = ? & \Rightarrow y(\tau_{\mathrm{c}}) = y(2\tau_{\mathrm{ta}}) = 0 \\ Q_{\mathrm{cta}}: & \rho_{\mathrm{de}} = ? & \Rightarrow \boxed{Assuming} \ \delta_{\mathrm{de,ta}}^{\mathrm{NL}} = r \delta_{\mathrm{m,ta}}^{\mathrm{NL}} \\ & \Rightarrow Q_{\mathrm{cta}} = \frac{Q_{\mathrm{ta}}}{1 + r[\zeta(w, z_{\mathrm{ta}}, r) - 1]} \\ ^*d\tau \equiv H_{\mathrm{ta}} \sqrt{\Omega_{\mathrm{m}}(x_{\mathrm{ta}})} dt \ , \ x \equiv \frac{a}{a_{\mathrm{ta}}}, \ y \equiv \frac{R}{R_{\mathrm{ta}}} \int_{\rho_{\mathrm{de}}} \left| \frac{\rho_{\mathrm{m}}}{\bar{\rho}_{\mathrm{m}}} \right|_{z_{\mathrm{ta}}}, \\ Q_{\mathrm{ta}} \equiv \frac{\bar{\rho}_{\mathrm{m}}}{\bar{\rho}_{\mathrm{de}}} \left| \frac{\Omega_{\mathrm{m}}^{0}}{2} \left(1 + z_{\mathrm{ta}}\right)^{-3w}, \ Q_{\mathrm{cta}} \equiv \frac{\bar{\rho}_{\mathrm{m}}}{\rho_{\mathrm{de}}} \right|_{z_{\mathrm{ta}}}, \\ \delta_{\mathrm{m}}^{\mathrm{NL}} \equiv \rho_{\mathrm{m}}/\bar{\rho}_{\mathrm{m}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} \equiv \rho_{\mathrm{de}}/\bar{\rho}_{\mathrm{de}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} \equiv \rho_{\mathrm{de}}/\bar{\rho}_{\mathrm{de}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} \equiv \rho_{\mathrm{de}}/\bar{\rho}_{\mathrm{de}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} \equiv \rho_{\mathrm{de}}/\bar{\rho}_{\mathrm{de}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} \equiv \rho_{\mathrm{de}}/\bar{\rho}_{\mathrm{de}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} \equiv \rho_{\mathrm{de}}/\bar{\rho}_{\mathrm{de}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} \equiv \rho_{\mathrm{de}}/\bar{\rho}_{\mathrm{de}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} \equiv \rho_{\mathrm{de}}/\bar{\rho}_{\mathrm{de}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} \equiv \rho_{\mathrm{de}}/\bar{\rho}_{\mathrm{de}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} \equiv \rho_{\mathrm{de}}/\bar{\rho}_{\mathrm{de}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} \equiv \rho_{\mathrm{de}}/\bar{\rho}_{\mathrm{de}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} \equiv \rho_{\mathrm{de}}/\bar{\rho}_{\mathrm{de}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} \equiv \rho_{\mathrm{de}}/\bar{\rho}_{\mathrm{de}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} \equiv \rho_{\mathrm{de}}/\bar{\rho}_{\mathrm{de}} - 1 \ , \ \text{and} \ \delta_{\mathrm{de}}^{\mathrm{NL}} = 0 \ , \ \ \delta_{\mathrm{de}}^{\mathrm{de}} = 0 \ , \ \ \delta_{\mathrm{de}}^{\mathrm{d$$

clustering DE ($\alpha = 0$), $\rho_{\rm de} \neq \bar{\rho}_{\rm de}$







- the upper: r = 1, w = -0.8, $z_{ta} = 0.8$
- $\Omega_{\rm m,0} = 0.274$
- * $r = \delta_{
 m de,ta}^{
 m NL}/\delta_{
 m m,ta}^{
 m NL}$

$*$
 $\zeta \equiv rac{
ho_{
m m}}{ar
ho_{
m m}}ig|_{z_{
m ta}}$

$\sf Virialization$

(Maor and Lahav, Astropart. Phys. 07 (2005) 003)

virial theorem and energy conservation: $\left[U + \frac{R}{2} \frac{\partial U}{\partial R}\right]_{\perp} = U_{\rm ta}$

$$U = \frac{1}{2} \int_0^R \rho_{\mathrm{m}} \Phi_{\mathrm{m}} dV + \frac{1}{2} \int_0^R \rho_{\mathrm{de}} \Phi_{\mathrm{m}} dV + \frac{1}{2} \int_0^R \rho_{\mathrm{m}} \Phi_{\mathrm{de}} dV + \frac{1}{2} \int_0^R \rho_{\mathrm{de}} \Phi_{\mathrm{de}} dV$$

whole system virializing

$$[1 + (2+3w)q + (1+3w)q^{2}]y_{\text{vir}} - \frac{1}{2}(2+3w)(1-3w)qy_{\text{vir}}^{-3w} - \frac{1}{2}(1+3w)(1-6w)q^{2}y_{\text{vir}}^{-6w} = \frac{1}{2}$$

only matter virializing

$$(1+q)y_{\rm vir} - \frac{q}{2}(1-3w)y_{\rm vir}^{-3w} = \frac{1}{2}$$

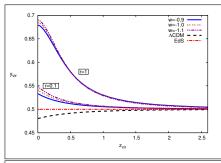
EdS universe

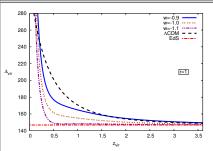
$$y_{
m vir}=rac{1}{2}$$

*
$$q \equiv \left. \frac{
ho_{
m de}}{
ho_{
m m}} \right|_{z_{
m ta}} = \left(\frac{\delta_{
m m,t}^{
m HL},{
m ta}+1}{\delta_{
m m,t}^{
m HL},{
m ta}+1} \right) \frac{1-\Omega_{
m m,0}}{\Omega_{
m m,0}} (1+z_{
m ta})^{3w} = \left(\frac{r(\zeta-1)+1}{\zeta} \right) \frac{1-\Omega_{
m m,0}}{\Omega_{
m m,0}} (1+z_{
m ta})^{3w} = \left(\frac{r(\zeta-1)+1}{\zeta} \right) \frac{1-\Omega_{
m m,0}}{\Omega_{
m m,0}} (1+z_{
m ta})^{3w}$$



Virialization



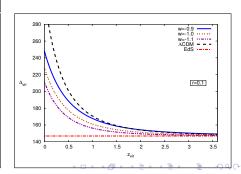


•
$$\Omega_{\rm m,0} = 0.274$$

*
$$y_{\rm vir} \equiv \frac{R_{\rm vir}}{R_{\rm ta}}$$

*
$$\Delta_{\text{vir}} \equiv \frac{\rho_{\text{m}}}{\bar{\rho}_{\text{m}}} \Big| = \zeta \left(\frac{\mathsf{x}_{\text{vir}}}{\mathsf{v}_{\text{vir}}}\right)^{\zeta}$$

*
$$r = \delta_{
m de,ta}^{
m NL}/\delta_{
m m,ta}^{
m NL}$$



continuity eq.:

$$\dot{\rho}_{\mathrm{x}} + \overrightarrow{\nabla} \cdot [(\rho_{\mathrm{x}} + P_{\mathrm{x}}) \vec{v}] = 0$$

Euler eq.:

$$\dot{\vec{v}} + (\vec{v} \cdot \overrightarrow{\nabla})\vec{v} = -\overrightarrow{\nabla}\Phi$$

adding small perturbations:

$$ho_{\mathrm{x}} = ar{
ho}_{\mathrm{x}}(1 + \delta_{\mathrm{x}}), \ \ \overrightarrow{\mathsf{V}} = H\overrightarrow{\mathsf{x}} + \overrightarrow{\mathsf{u}}$$

Poisson eq.:

$$\overrightarrow{\nabla}^2 \delta \Phi = 4\pi G \left(\delta \rho_{\mathrm{m}} + \delta \rho_{\mathrm{de}} + 3\delta P_{\mathrm{de}} \right) a^2$$



 δ in Linear Regime

(*
$$\tilde{x}\equiv a/a_0$$
, $\eta\equiv\sqrt{\Omega_{\mathrm{m},0}}H_0t$)

background universe

$$rac{d ilde{x}}{d\eta} = rac{1}{\sqrt{ ilde{x}\Omega_{
m m}(ilde{x})}} \qquad ext{with } \Omega_{
m m}(ilde{x}) = \left(1 + rac{1 - \Omega_{
m m}^0}{\Omega_{
m m}^0} ilde{x}^{-3w}
ight)^{-1}$$

matter dominated universe: $\tilde{x}_i = (3\eta_i/2)^{2/3}$

density perturbation

$$egin{aligned} rac{d^2\delta_{
m m}}{d\eta^2} + rac{2}{ ilde{x}}rac{d ilde{x}}{d\eta}rac{d\delta_{
m m}}{d\eta} &= rac{3}{2 ilde{x}^3}\left[\delta_{
m m} + rac{1-\Omega_{
m m}^0}{\Omega_{
m m}^0}(1+3w) ilde{x}^{-3w}\delta_{
m de}
ight] \ rac{d\delta_{
m de}}{d\eta} &= (1+w)rac{d\delta_{
m m}}{d\eta} \Rightarrow \delta_{
m de} &= (1+w)\delta_{
m m} \end{aligned}$$

matter dominated universe:

$$\delta_{
m m} \propto a \Rightarrow d\delta_{
m m}/d\eta|_i = 2\delta_{
m m,i}/(3\eta_i)$$



 δ in Non-Linear Regime (Creminelli et al., JCAP 03 (2010) 027)

- ullet background universe $rac{d ilde{x}}{d\eta} = rac{1}{\sqrt{ ilde{x}\Omega_{
 m m}(ilde{x})}}$
- spherical overdense region $\sqrt{x_1 x_m(x_2)}$

$$\frac{d^2\tilde{y}}{d\eta^2} + \frac{1}{2} \left[\frac{1 + \delta_{\text{m},i}}{\tilde{x}_i^3} \frac{1}{\tilde{y}^2} + (1 + 3w)(1 + \delta_{\text{de}}^{\text{NL}}) \frac{1 - \Omega_{\text{m}}^0}{\Omega_{\text{m}}^0} \frac{\tilde{y}}{\tilde{x}^{3(1+w)}} \right] = 0$$

 $\tilde{y}(\eta_i) = \tilde{y}_i = 1$ for $\tilde{y} \equiv \frac{R}{R_i}$ • density contrast

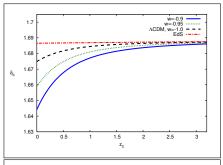
$$\frac{d\delta_{\mathrm{de}}^{\mathrm{NL}}}{d\eta} + 3(1+w)(1+\delta_{\mathrm{de}}^{\mathrm{NL}})\left(\frac{1}{\tilde{y}}\frac{d\tilde{y}}{d\eta} - \frac{1}{\tilde{x}}\frac{d\tilde{x}}{d\eta}\right) = 0$$

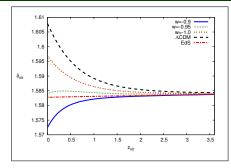
linear in early time: $\delta^{\rm NL}_{{
m de},i} pprox \delta_{{
m de},i} = (1+w)\delta_{{
m m},i}$

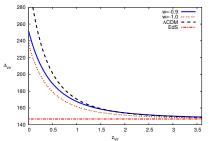
$$rac{d\delta_{
m m}^{
m NL}}{dn} + 3(1+\delta_{
m m}^{
m NL})\left(rac{1}{ ilde{arkappa}}rac{d ilde{y}}{dn} - rac{1}{ ilde{x}}rac{d ilde{x}}{dn}
ight) = 0$$

linear in early time:

$$|d ilde{y}/d\eta|_ipprox 2(1-\delta_{\mathrm{m},i}/3)/(3\eta_i)| \ \ (\delta_{\mathrm{m}}\propto ilde{x}\propto \eta^{2/3})$$







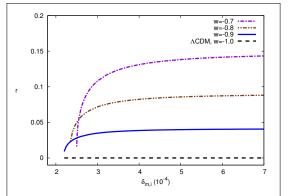
•
$$\Omega_{\rm m,0} = 0.274$$

*
$$\delta_{\rm c} = \delta_{\rm m}(\eta_{\rm c})$$

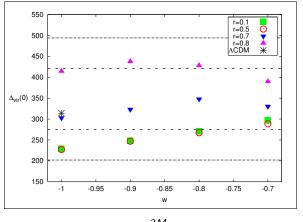
*
$$\delta_{\rm vir} = \delta_{\rm m}(\eta_{\rm vir})$$

*
$$\Delta_{
m vir} = \left(1 + \delta_{
m m,ta}^{
m NL}
ight) \left(rac{a_{
m vir}/a_{
m ta}}{R_{
m vir}/R_{
m ta}}
ight)^3$$

	1	II
$ ho_{ m m}$	introduce $\zeta\equiv rac{ ho_{ m m}}{ar ho_{ m m}}igg _{z_{ m ta}}$ satisfying $y(au_{ m c})=y(2 au_{ m ta})=0$	give a large enough $\delta_{\mathrm{m},i}$
$\rho_{ m de}$	introduce r with $\delta_{ m de,ta}^{ m NL} = r \delta_{ m m,ta}^{ m NL}$	$\delta_{{ m de},i}^{ m NL}pprox \delta_{{ m de},i}=(1+{\it w})\delta_{{ m m},i}$



(Crook et al., ApJ 655 (2007) 790)

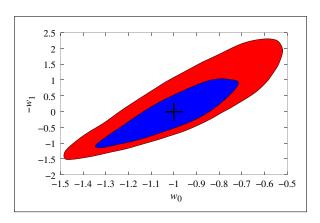


$$\Delta_{
m vir} = rac{3 \emph{M}_{
m vir}}{4 \pi
ho_{
m c,0} \Omega_{
m m,0} (1+z)^3 \emph{R}_{
m vir}^3}$$

• It is likely that $0.5 < r \le 0.8$ for models containing clustered DE with $w \le -0.9$ at a confidence level of 68%.

CPL (Chevallier-Polarski-Linder) model:

$$w(a) = w_0 + w_1(1-a)$$



(U. Seljak et al., PRD 71, 103515 (2005))



CPL (Chevallier-Polarski-Linder) model:

$$egin{array}{lcl} w(a) &=& w_0 + w_1(1-a) \ &=& \left[w_0 + w_1(1-a_{
m ta})
ight] + w_1 a_{
m ta} \left(1 - rac{a}{a_{
m ta}}
ight) \ &=& w_{
m ta} + w_{
m 1,ta} \left(1 - rac{a}{a_{
m ta}}
ight) \end{array}$$

$$w_b\left(rac{a}{a_{ ext{ta}}}
ight) = w_{ ext{ta}} + w_{1, ext{ta}}\left(1 - rac{a}{a_{ ext{ta}}}
ight) \qquad w_{ ext{ta}} = w_0 + w_1(1 - a_{ ext{ta}}) \ = w_0 + rac{w_1 z_{ ext{ta}}}{1 + z_{ ext{ta}}} \ w_{1, ext{ta}} = w_1 a_{ ext{ta}} = rac{w_1}{1 + z_{ ext{ta}}}$$

In this way, $w_c = w_b = w_{\rm ta}$ at $t = t_{\rm ta}$.

Background Universe

$$\begin{split} \frac{\ddot{a}}{a} &= -\frac{4\pi G}{3} \left[\bar{\rho}_{\mathrm{m}} + (1 + 3w_b) \bar{\rho}_{\mathrm{de}} \right] \\ \dot{\bar{\rho}}_{\mathrm{m}} &+ 3 \left(\frac{\dot{a}}{a} \right) \bar{\rho}_{\mathrm{m}} = 0 \\ \dot{\bar{\rho}}_{\mathrm{de}} &+ 3(1 + w_b) \left(\frac{\dot{a}}{a} \right) \bar{\rho}_{\mathrm{de}} = 0 \end{split}$$

Spherical Overdensity

$$egin{align} rac{\ddot{R}}{R} &= -rac{4\pi G}{3} \left[
ho_{
m m} + (1+3w_c)
ho_{
m de}
ight] \ \dot{
ho}_{
m m} + 3\left(rac{\dot{R}}{R}
ight)
ho_{
m m} = 0 \ \dot{
ho}_{
m de} + 3(1+w_c)\left(rac{\dot{R}}{R}
ight)
ho_{
m de} = 0 \ \end{aligned}$$

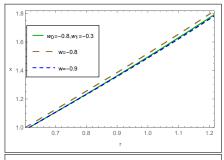
$$\Rightarrow \frac{\bar{\rho}_{de}}{\bar{\rho}_{de,ta}} = \exp\left[3\int_{a/a_{ta}}^{1} \frac{1+w_{b}(u)}{u} du\right] = f\left(\frac{a}{a_{ta}}\right)$$

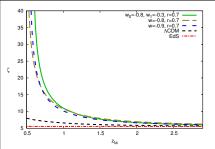
$$= x^{-3(1+w_{0}+w_{1})} e^{-3w_{1}(1-x)/(1+z_{ta})}$$

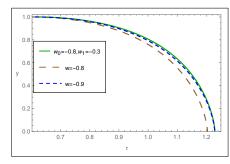
$$\Rightarrow \frac{\rho_{de}}{\rho_{de,ta}} = \exp\left[3\int_{R/R_{ta}}^{1} \frac{1+w_{c}(u)}{u} du\right] = f\left(\frac{R}{R_{ta}}\right)$$

$$= y^{-3(1+w_{0}+w_{1})} e^{-3w_{1}(1-y)/(1+z_{ta})}$$

$$\frac{dx}{d\tau} = \sqrt{x^{-1} + \frac{1}{Q_{ta}} f\left(\frac{a}{a_{ta}}\right) x^{2}} \\
= \sqrt{x^{-1} + \frac{1}{Q_{ta}} x^{-1 - 3(w_{0} + w_{1})} e^{-3w_{1}(1-x)/(1+z_{ta})}} \\
\Rightarrow \tau_{ta} = \int_{0}^{1} \frac{dx'}{\sqrt{x'^{-1} + \frac{1}{Q_{ta}} x'^{-1 - 3(w_{0} + w_{1})} e^{-3w_{1}(1-x')/(1+z_{ta})}}} \\
\frac{d^{2}y}{d\tau^{2}} = -\frac{1}{2} \left[\zeta y^{-2} + \left(\frac{1+3w_{c}}{Q_{cta}}\right) f\left(\frac{R}{R_{ta}}\right) y \right] \\
= -\frac{1}{2} \left\{ \zeta y^{-2} + \frac{1}{Q_{cta}} \left[1+3w_{0} + 3w_{1} \left(1-\frac{y}{1+z_{ta}}\right) \right] \right. \\
y^{-2-3(w_{0} + w_{1})} e^{-3w_{1}(1-y)/(1+z_{ta})} \right\} \\
\star Q_{ta} \equiv \frac{\tilde{\rho}_{m}}{\tilde{\rho}_{de}} \Big|_{z_{ta}} = \frac{Q_{ta}}{1+r[\zeta(w_{0}, w_{1}, z_{ta}, r) - \frac{1}{2}]} \\
\star Q_{cta} \equiv \frac{\tilde{\rho}_{m}}{\tilde{\rho}_{de}} \Big|_{z_{ta}} = \frac{Q_{ta}}{1+r[\zeta(w_{0}, w_{1}, z_{ta}, r) - \frac{1}{2}]} \\
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\star Q_{cta} \equiv \frac{\tilde{\rho}_{m}}{\tilde{\rho}_{de}} \Big|_{z_{ta}} = \frac{Q_{ta}}{1+r[\zeta(w_{0}, w_{1}, z_{ta}, r) - \frac{1}{2}} \Big|_{z_{ta}} = \frac{\tilde{\rho}_{ct}}{\tilde{\rho}_{ct}} \Big|_{z_{ta}} = \frac{\tilde{\rho}_{ct}}{1+r[\zeta(w_{0}, w_{1}, z_{ta}, r) - \frac{1}{2}} \Big|_{z_{ta}} = \frac{\tilde{\rho}_{c$$







•
$$\Omega_{\rm m,0} = 0.274$$

•
$$r = 0.7$$

*
$$X \equiv \frac{a}{a_{\rm ta}}$$

*
$$y \equiv \frac{R}{R_{ta}}$$

*
$$r = \delta_{\rm de,ta}^{\rm NL}/\delta_{\rm m,ta}^{\rm NL}$$

*
$$\zeta \equiv \frac{\rho_{\rm m}}{\bar{\rho}_{\rm m}}$$



Conclusions

- We presuppose $\delta_{
 m de,ta}^{
 m NL}=r\delta_{
 m m,ta}^{
 m NL}$ to remove the difficulty with clustered DE in the SCM.
- We treat the overdense region as an isolate system and determin $y_{\rm vir} (=R_{\rm vir}/R_{\rm ta})$ and the nonlinear overdensity $\Delta_{\rm vir}$. Also we distinguish $\Lambda {\rm CDM}$ from the clustering model of DE with w=-1.
- $\Delta_{\rm vir}$ in $\Lambda{\rm CDM}$ reaches the maximum value while for other clustered DE models the amplitude of $\Delta_{\rm vir}$ is proportional to w.
- We track down the complete evolution of the clustered matter and DE by combining SCM with the linear evolution of density perturbation at early times, equivalent to introducing the new parameter r.
- We have found that the criterion $\delta_{\mathrm{m},i} \geq 2.2 \times 10^{-4} \sim 2.5 \times 10^{-4}$ depending on w has to be fulfilled to form a cosmic structure.
- The ratio r is less dependent on the initial matter density contrast for cases with $\delta_{\mathrm{m},i} \geq 5 \times 10^{-4}$.
- We use observational data of galaxy clusters to constrain the parameter r.
- Our method can be extended to DE with time-varying w following from the CPL model.

