

Towards the Theory of

Quantum Relativity

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人之能知天也，
能知人之天也。

~王夫之

Symmetry of “ **Space -Time-U...**” : -
the New Arena for Fundamental Physics

Galilean Relativity

Einstein Relativity — special and general

(Special) Quantum Relativity

or Doubly (Triply) / Deformed Special Relativity

H.Synder / C.N.Yang (1947)

G.Amelino-Camelia (2001)

L.Smolin + (2002)

..... General Quantum Relativity \Rightarrow **Quantum Gravity +**

My Basic Perspectives :-

— on issues related to Quantum Space-Time

- need new understanding of “Space-Time” (relativity principle)

- within Space-Time Geometry

Non-Commutative Geometry *is to* Quantum Gravity
as Non-Euclidean Geometry *is to* (Einstein) Gravity

- Quantum Field Theory, on (Minkowski/Einstein) Space-Time, perspective would be valid way beyond the experimentally probed energy/length scale — *a faith* I don't share (*Penrose*)

— Linear (Vs Nonlinear) Realization

- an exploration into **new** ways to think about fundamental physics
- withhold phenomenological discussions

A Stable Symmetry :-

R.V. Mendes (1994)
Chryssomalakos & Okon (2004)

- perturbations (deformations) of a symmetry algebra
- stable if deformed algebras isomorphic to original
- structure constants **perturbation Vs experimental uncertainty**
- Examples —
 - 1/ Galilean to **Einstein Relativity** ($1/c^2 \neq 0$)
 - 2/ classical to **quantum mechanics** ($\hbar \neq 0$)

Poisson bracket \implies Moyal bracket
- **deformation parameters** \longrightarrow **fundamental constants**

Space-Time + Quantum Symmetry

‘Poincaré + Heisenberg’ algebra

\implies **SO(1,5)** OR **SO(2,4)**

WI Contraction :-

$$SO(1, 3) \longrightarrow ISO(3)$$

$$[N_i, N_j] = -i \frac{1}{c^2} M_{ij}$$

$\frac{1}{c^2} \rightarrow 0$ boosts commute (i.e. Galilean)

- $SO(1, 3)$ as the stablization :

$$[M_{ij}, L_{0k}] = i (\eta_{jk} L_{0i} - \eta_{ik} L_{0j}) \qquad [L_{0i}, L_{0j}] = -i M_{ij}$$

take $N_i \rightarrow \frac{1}{c} L_{0i}$: isomorphic

? value of c : 3×10^8 Vs 1 or 3×10^{-7} (km ps⁻¹)

Quantum Relativity $SO(1, 5)$ or $SO(2, 4)$:-

$$[J_{AB}, J_{LN}] = i (\eta_{BL} J_{AN} - \eta_{AL} J_{BN} + \eta_{AN} J_{BL} - \eta_{BN} J_{AL})$$

with $\eta^{AB} = (1, -1, -1, -1, -1, \mp 1)$.

- linear realization on 6-geometry — $J_{AB} = i (x_A \partial_B - x_B \partial_A)$
- J_{AB} for $0 \rightarrow 3$ are the 10 generator for Lorentz symmetry $M_{\mu\nu}$
- NOT necessarily means extra space-time dimensions
- 3 deformations — $ISO(3) \rightarrow SO(1, 3) \hookrightarrow ISO(1, 3) \rightarrow SO(1, 4) \hookrightarrow ISO(1, 4) \rightarrow SO(1, 5)$ or $SO(2, 4)$
- 3 invariant quantities : c κ ℓ Kowalski-Glikman & Smolin(2004)
OK (2007)
- curved momentum space \Rightarrow non-commutative space-time
+ curved space-time \Rightarrow non-commutative momentum space

The Three Deformations :-

$$SO(1, 4) :- |z^A| \leq \ell$$

$$\eta_{AB} z^A z^B = -\ell^2 \left(1 + \frac{1}{G^2}\right)$$

$\Delta x^i(t) = v^i \cdot t$	$\Delta x^\mu(\sigma) = p^\mu \cdot \sigma$	$\Delta x^A(\rho) = z^A \cdot \rho$
$ v^i \leq c$ $-\eta_{ij} v^i v^j = c^2 \left(1 - \frac{1}{\gamma^2}\right)$	$ p^\mu \leq \kappa c$ $\eta_{\mu\nu} p^\mu p^\nu = \kappa^2 c^2 \left(1 - \frac{1}{\Gamma^2}\right)$	$ z^A \leq \ell$ $\eta_{AB} z^A z^B = \ell^2 \left(1 - \frac{1}{G^2}\right)$
$M_{0i} \equiv N_i \sim P_i$ $[N_i, N_j] \longrightarrow -i M_{ij}$	$J_{\mu 4} \equiv O_\mu \sim P_\mu$ $[O_\mu, O_\nu] \longrightarrow i M_{\mu\nu}$	$J_{A5} \equiv O'_A \sim P_A$ $[O'_A, O'_B] \longrightarrow i J_{AB}$
$\vec{u}^4 = \frac{\gamma}{c}(c, v^i)$ $\eta_{\mu\nu} u^\mu u^\nu = 1$ $\mathbb{R}^3 \rightarrow SO(1, 3)/SO(3)$	$\vec{\pi}^5 = \frac{\Gamma}{\kappa c}(p^\mu, \kappa c)$ $\eta_{AB} \pi^A \pi^B = -1$ $M^4 \rightarrow SO(1, 4)/SO(1, 3)$	$\vec{X}^6 = \frac{G}{\ell}(z^A, \ell)$ $\eta_{\mathcal{M}\mathcal{N}} X^\mathcal{M} X^\mathcal{N} = -1$ $M^5 \rightarrow SO(1, 5)/SO(1, 4)$

- 4-momentum is defined by $p^\mu = \frac{dx^\mu}{d\sigma}$
- $z^A = \frac{dx^\mu}{d\rho}$ is chosen as a length
- without the p^μ deformation, $X^\mathcal{M}$ are just alternative coordinates of a de-Sitter space-time
- no $\hookrightarrow ISO(1, 5)/ISO(2, 4)$: no naive *translation* symmetries

Momentum Boosts (with \hat{X}_μ) :-

$$J_{\mu 4} = i (x_\mu \partial_4 - x_4 \partial_\mu)$$

— rotations among x_μ 's and $x_4 = -\kappa c \sigma$

A.Das & O.K.

★ σ is peculiar — (*doubt* space-time interpretation !)

• $ISO(1, 3) \rightarrow SO(1, 4)$ since Snyder

— κ as Planck mass

• preferred frame for a quantum state : $\vec{\pi}^5 = (0, 0, 0, 0, 1)$

the reference frame does not see its own 4-momentum/mass

• boost characterized by p^μ , hence mass dependent

— quantum frames

- $p^\mu = \frac{dx^\mu}{d\sigma}$ Vs $m c u^\mu = \gamma m c \frac{dx^\mu}{dx^0}$ (as Einstein/classical limit)
 $\sigma \longrightarrow \frac{\tau}{m}$ for an Einstein particle
 recall : “space-like” geometrical signature
- general transformation — on-shell condition not preserved
 classical to quantum frame, uncertainty : observer to observed
- σ -coordinate Vs literature
 - Vs proper time / covariant Schrödinger QM
 - system with indefinite mass
 - QM on M^4 stochastic metric fluctuation
- ★ momentum boosts needed for group quantization (relativistic)

Quantum Frame of Reference :-

Aharonov & Kaufherr, PR 30, 368
C. Rovelli, Class.Quantum.Grav. 8, 317

- context: nonrelativistic QM / gravitation
- characterized by *mass* (together with $v + \dots$)
- interaction of frame (observer/device) with observed
— observation has nontrivial effect on observer
- uncertainty principle on frame (observer/device)
—→ uncertainty on observed
- grav. — local gauge inv. observables obtained after dynamic properties of frame (observer/device) taken into consideration

★ *Is there such a thing as a practical quantum measurement ?*

Translational Boosts (with $P_a^{(\mathcal{L})}$ — \hat{P}_μ) :-

$$J_{A5} = i (x_A \partial_5 - x_5 \partial_A)$$

— rotations among x_A 's and $x_5 = -\ell \rho$

A.Das & O.K.
O.K.

★ ρ may characterize the scale

- boost characterized by z^A , a location vector
sort of a translation

- preferred frame for a state : $\vec{X}^6 = (0, 0, 0, 0, 0, 1)$

the reference frame does not see its own location

— $(0, 0, 0, 0, 0, 1)$ characterizes the coordinate origin $z^A = 0$

★ a reference frame does not observe itself (no v^i , p^μ , $z^A = 0$)

- Beltrami description *preserves* physics of 5-momentum constraint

- $ISO(1, 4) \rightarrow SO(1, 5)$

A.Das & O.K.

- ℓ an IF cutoff (Vs κ) \leftrightarrow cosmological constant

Kowalski-Glikman & Smolin(2004)

- dS_5 as “quantum” world

- $ISO(1, 4) \rightarrow SO(2, 4)$

O.K.

- ℓ_P as Planck length (UV cutoff); $\kappa c \ell_P = \hbar$

$$(\kappa^2 = c \hbar / G, \quad \ell_P^2 = \hbar G / c^3)$$

- AdS_5 as quantum world

- no $ISO(1, 5)$ or $ISO(2, 4)$ — no coordinate translations

- quantum general relativity — the role of G ?

Non-commutative (Space-Time) Geometry :-

cf. Kowalski-Glikman & Smolin
Chryssomalakos & Okon

- (four) space-time position operators $\hat{X}_\mu = -\frac{1}{\kappa c} i (x_\mu \partial_4 - x_4 \partial_\mu)$

$$[\hat{X}_\mu, \hat{X}_\nu] = \frac{i}{\kappa^2 c^2} M_{\mu\nu}$$

- (four) energy-momentum operators $\hat{P}_\mu = \frac{1}{\ell} i (x_\mu \partial_5 - x_5 \partial_\mu)$

$$[\hat{P}_\mu, \hat{P}_\nu] = \frac{i}{\ell^2} M_{\mu\nu}$$

$$[\hat{X}_\mu, \hat{P}_\nu] = -i \eta_{\mu\nu} \hat{F}, \quad [\hat{X}_\mu, \hat{F}] = +\frac{i}{\kappa^2 c^2} \hat{P}_\mu, \quad [\hat{P}_\mu, \hat{F}] = -\frac{i}{\ell^2} \hat{X}_\mu$$

- $i \partial_\mu \ll \kappa c$ and $i \partial_4 = p_4 = -\kappa c$: $\hat{X}_\mu \longrightarrow x_\mu$

- $x_\mu \ll \ell$ and $x_5 = -\ell$ ($\rho = 1$) : $\hat{P}_\mu \longrightarrow i \partial_\mu = p_\mu$

$SO(2, 4)$ Quantum Relativity :-

— AdS_5 as quantum world

- matching to 4D conformal symmetry :-

$$K_\mu \Rightarrow \sqrt{2} J_{\mu-} \quad P_\mu \Rightarrow \sqrt{2} J_{\mu+} \quad D \Rightarrow -J_{45}$$

$$J_{\mu\pm} \equiv i\hbar (x_\mu \partial_\pm - x_\pm \partial_\mu) = \frac{1}{\sqrt{2}} (J_{\mu 5} \pm J_{\mu 4})$$

- coordinate transformation — (gives the right matching)

$$(x^\mu, x^4, x^5) \longrightarrow (y^\mu, \frac{1}{2} \eta_{\mu\nu} y^\mu y^\nu + \frac{1}{2}, \frac{1}{2} \eta_{\mu\nu} y^\mu y^\nu - \frac{1}{2})$$

- However : $\Rightarrow x_+ = x^- = \frac{1}{\sqrt{2}} (x^5 - x^4) = -\frac{1}{\sqrt{2}}$

$$x_- = x^+ = \frac{1}{\sqrt{2}} y^2 \quad \partial_+ = 0 \quad \partial_5 = -\partial_4 = \frac{1}{\sqrt{2}} \partial_- = x^\nu \partial_\nu$$

★ $\eta_{\mathcal{M}\mathcal{N}} x^\mathcal{M} x^\mathcal{N} = \eta_{\mu\nu} x^\mu x^\nu + 2x_+ x_- = 0$ ($\neq -1$) for conformal universe

★ x^5, x^4 (or $x_- = x^+$) translations related to scaling ?

Remarks :-

- (*bottom line*) an interesting, radical but sensible, approach
- primitive stage — difficult to make (and identify) progress
- need creative but careful thinking about ‘physics’
 beyond usual framework
- job for Einstein — hope we can make minor steps
 “ The chief cause of my failure was my clinging to the idea that the
 variable t only can be considered as the true time and my local t' must
 In Einstein’s theory, t' plays the same part as t ”

The Theory of Electrons (1916 ed.) — Lorentz (from A.I. Miller)

Space-Time-U ...
as the New Arena for Fundamental Physics

THANK YOU !