

Dynamics of Spacetime from a Large N Matrix Model

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Dynamics of Spacetime from a Large N Matrix Model

Yang-Mills (YM) type: $-\text{tr}[A,A][A,A]/4$

1. Introduction and Motivation

String Theory

- *Fundamental degrees of freedom are extended object, not point like particles*



- *A nice candidate for the unified quantum theory with gravity*
- *To investigate strong gravity phenomenon (Black Hole, very early Universe), 2nd quantized formulation may be needed.*

Main remarkable properties from this extension:

1. Infinitely many state of string vibration



*Infinitely many one particle state
(include graviton, gauge boson, ...)*

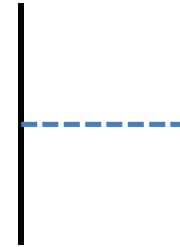
***2. Formulated in higher dimension
(e. g. $D=26$, $D=10$ flat space)***

3. UV finiteness of loop amplitude

String theory is a nice candidate of the unified quantum theory including gravity.

- *Well established formulation is 1-st quantized/
perturbative formulation*

*Coupling constant of gravitational correction
in perturbation theory ($D=10$):*



$$(E^8 G_N)$$

G_N : Newton constant in 10D

Perturbation theory will break down in processes with very high energy.

There will be large back reaction to spacetime if

$$E > \frac{1}{g^{1/4} \sqrt{\alpha'}}$$

g : string coupling. α' : (string length scale)².

***There could be black hole production
with Schwatzschild radius:***

$$R \sim (G_N E)^{1/7}$$

- ***To investigate strong gravity phenomenon (Black Hole, very early Universe), 2nd quantized formulation may be needed.***

Proposals for 2nd quantized theory

- *String Field Theory*

Reproduces perturbative expansion of S-matrix

Still difficult to see non-perturbative effect

(Recent Progresses :Tachyon condensation, D-brane solution,)

- ✓ ▪ *Reformulation by Large N Matrix Theory*

Proposals for critical superstring in recent 15 years

1 dim model (BFSS), 0 dim model (IKKT), etc.

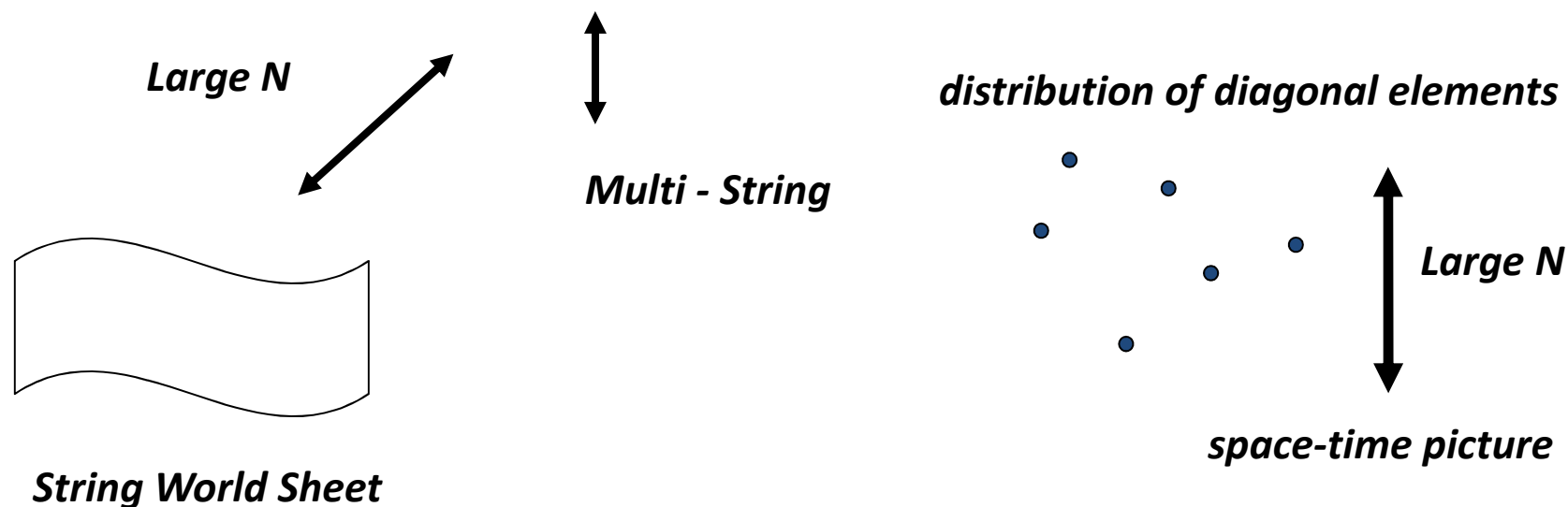
Matrix Theory proposal (I) : IKKT model

Ishibashi, Kawai, Kitazawa, Tsuchiya (1996)

Fundamental degrees of freedom:
 $N \times N$ hermitian matrices

$A_{(a)}, \psi$

$$S = -\frac{1}{4g^2} \text{tr} \left([A_a, A_b][A^a, A^b] \right) + \frac{1}{2g^2} \text{tr} \left(\bar{\psi} \Gamma^a [A_a, \psi] \right)$$



Matrix Theory proposal (II) : BFSS model

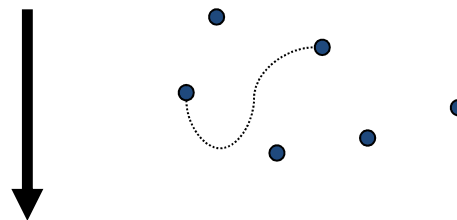
Banks, Fishler, Shenker, Susskind (1996)

Large N matrix quantum mechanics

$$X^{(I)}(t), A_{(a)}(t), \Psi(t)$$

$$S_{BFSS} = \int dt \operatorname{tr} \left(-\frac{1}{2} [D_0, X_{(i)}] [D^0, X^{(i)}] + \frac{1}{4} [X_{(i)}, X_{(j)}] [X^{(i)}, X^{(j)}] - \frac{1}{2} \bar{\Psi} D_0 \Psi - \frac{1}{2} \bar{\Psi} \Gamma^{(i)} [X^{(i)}, \Psi] \right)$$

D0 brane effective theory in type IIA superstring



Large N limit formulates M-theory on a circle in light-cone direction

Having these Yang-Mills matrix models, we have a possibility to study the dynamics of spacetime that emerges from matrix.

→ *“Quantum theory of gravity”*

Purpose of this talk:

*To introduce **a part of** studies that have been done along this interest.*

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1. Introduction and Motivtion

***2. Matrix Description of
Space(time)***

Non-commutative spaces

Dynamics

***3. More on the Matrix
Description of Spacetime***

General curved spacetime

Dynamics

4. Summary and Outlook

Matrix Description of space(-time)

Space time may be described by matrix degrees of freedom

Dynamics of diagonal element

--> similar to the statistical system of polymer

There are toy examples that is more similar to ordinary space.

- *Two Examples*

1. *Fuzzy Sphere*
2. *Fuzzy Torus*

Matrix Description of space(-time) I: Fuzzy Sphere

Spin 1/2 operators are gives by 2×2 matrices (Pauli matrices)

$$J_1 = \begin{pmatrix} 0 & \frac{1}{2} \\ \frac{1}{2} & 0 \end{pmatrix}, \quad J_2 = \begin{pmatrix} 0 & -\frac{i}{2} \\ \frac{i}{2} & 0 \end{pmatrix}, \quad J_3 = \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & -\frac{1}{2} \end{pmatrix}$$

$$[J_a, J_b] = i\epsilon_{abc}J_c$$

Eigen states: $|l, m\rangle$

$$J_3|l, m\rangle = m|l, m\rangle$$

$$J^2|l, m\rangle = l(l+1)|l, m\rangle, \quad J^2 = J_1^2 + J_2^2 + J_3^2$$

For 2×2 matrices, $l=1/2$ and $m= 1/2, -1/2$

*Matrices J_a can be identified
with coordinate X_a (operator!)*

$|l, m\rangle$:

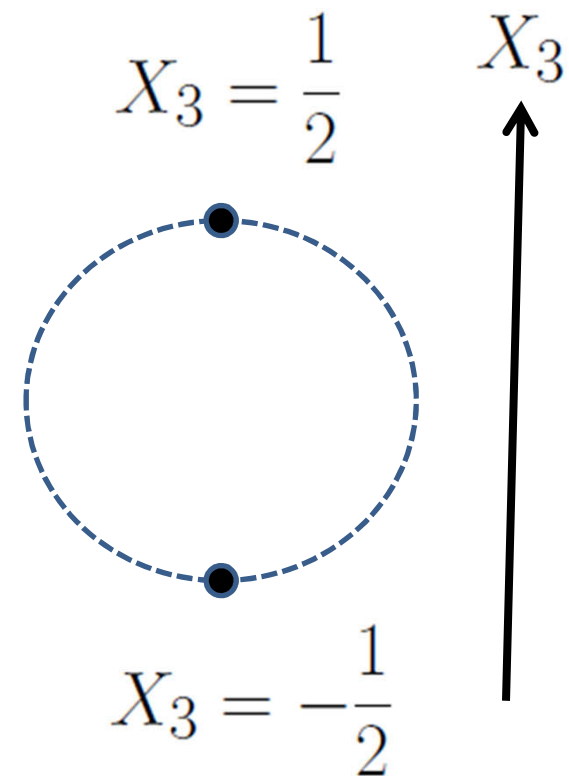
Eigenstate of coordinate operators

$$X_1^2 + X_2^2 + X_3^2 = \frac{3}{4}$$

This is a sphere with only 2-point.

$$[X_a, X_b] = i\epsilon_{abc}X_c$$

Coordinates are non-commutative \rightarrow “Fuzzy”



Field on Fuzzy Sphere

Hermitian Matrix

$$\begin{aligned} A &= \begin{pmatrix} a+d & b-ic \\ b+ic & a-d \end{pmatrix} = a \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + b \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + c \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + d \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\ &= 2aY_{0,0} + \sqrt{2}(b-ic)Y_{1,-1} + \sqrt{2}(b+ic)Y_{1,1} + 2dY_{1,0} \end{aligned}$$

$Y_{l,m}$: ***Matrix spherical harmonics***
Eigen matrix of

$$[J_3, Y_{l,m}] = mY_{l,m},$$

$$[J_1, [J_1, Y_{l,m}]] + [J_2, [J_2, Y_{l,m}]] + [J_3, [J_3, Y_{l,m}]] = l(l+1)Y_{l,m}.$$

$$\longleftrightarrow \quad \phi(\theta, \phi) = \sum_{l,m} \phi_{l,m} Y_{l,m}(\theta, \phi) \quad \text{on sphere.}$$

2×2 matrix $\rightarrow N \times N$ matrix

$$(N = 2l+1)$$

$$(Y_{p,m})_{s,t} = (-1)^{l-s} \sqrt{2p+1} \begin{pmatrix} l & p & l \\ -s & m & t \end{pmatrix}$$

$$p = 0, 1, \dots, 2l$$

3j-symbol

$$m = -p, -p+1, \dots, -1, 0, 1, \dots, p-1, p$$

$N \times N$ matrix $\rightarrow \infty \times \infty$ matrix ($N \rightarrow \infty$)

Continuum limit

Matrix description of space(time) II: Fuzzy Torus

$$U = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad V = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

property: $U^2 = V^2 = \mathbf{1}, \quad UV = (-1)VU$

Hermitian matrix

$$A = \begin{pmatrix} a + d & b - ic \\ b + ic & a - d \end{pmatrix} = a\mathbf{1} + dU + bV + icUV.$$

$$\longleftrightarrow \phi(\sigma_1, \sigma_2) = \sum_{m,n=0,1} \phi_{m,n} e^{im\sigma_1 + in\sigma_2}$$

$$U = e^{i\hat{\sigma}_1}, \quad V = e^{i\hat{\sigma}_2}$$

$$[\hat{\sigma}_1, \hat{\sigma}_2] = -i\pi \quad : \textit{non-commutativity}$$

Derivative:

$$[\hat{\sigma}_1, \] \Leftrightarrow -i\pi \frac{\partial}{\partial \sigma_2}, \quad [\hat{\sigma}_2, \] \Leftrightarrow -i\pi \frac{\partial}{\partial \sigma_1}.$$

Kinetic term:

$$\frac{1}{\pi^2} \sum_{i=1,2} [\sigma_i, \phi][\sigma_i, \phi] \Leftrightarrow \sum_{i=1,2} \partial_i \phi \partial_i \phi$$

2 × 2 matrix → N × N matrix

$$U_{ij} = \omega^{i-1} \delta_{ij}, \quad (\omega = e^{\frac{2\pi i}{N}}),$$

$$V_{ij} = \delta_{i,j-1},$$

$$U^N = 1, \quad V^N = 1, \quad UV = \omega^{-1} VU,$$

Non-commutativity:

$$[\hat{\sigma}_1, \hat{\sigma}_2] = -\frac{2\pi i}{N}$$

Matrix Model and Fuzzy Space (Dynamics)

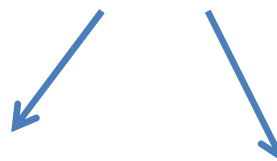
Decompose matrix A as $A=X+a$

X : back ground , a : field

Yang-Mills matrix model \rightarrow Gauge theory on X

$$-\frac{1}{4}\text{Tr}[A_i, A_j]^2$$

derivative



$$[A_1, A_2] = [X_1, X_2] + \underbrace{[X_1, a_1] - [X_2, a_1] + [a_1, a_2]}_{\text{derivative}}$$



$$(F_{12} = \partial_1 a_2 - \partial_2 a_1 + [a_1, a_2])$$

Quantum Correction

partition function: $Z = \int \mathcal{D}A e^{-S[A]}$

(A=X+a, integrate over a) $= \int \mathcal{D}X e^{-\Gamma[X]}$

Integration of X \approx Taking saddle point:
 $\delta\Gamma[X]/\delta X=0$

***Quantum correction of fuzzy space X
can be calculated in this way.***

Stability Analysis

- *fuzzy S^2 (2d)* Imai, Kitazawa, Takayama, D.T [03]
- *fuzzy $S^2 \times S^2$ (4d)* Imai, Kitazawa, Takayama, D.T [03]
Imai, Takayama [03]
- *fuzzy $S^2 \times S^2 \times S^2$ (6d)* Kaneko, Kitazawa, D.T [05]
- *fuzzy spaces with $SU(3)$ isometry (4d, 6d)* Kaneko, Kitazawa, , D.T [05]

4d fuzzy spaces are more stable than 2d, 6d fuzzy spaces.

More on the matrix description of space-time

We saw

Matrix can describe

***Non-commutative sphere, torus
Fields on these spaces***

***More general curved spacetime ?
(Black Hole, RW spacetime, etc.)***

Einstein Equation?

Approach by Hanada, Kawai, Kimura (HKK) [05]

Matrix with ∞ size



*Differential operator in
curved spacetime*

$$A_{(a)} = R_{(a)}{}^a(g) D_a(\nabla, \omega)$$

∇ : *covariant derivative*

ω : *spin connection*

Example: Covariant derivatives on 2-sphere

Two sphere (with stereographic coordinate)

$$\begin{aligned}\nabla_{(+)}^{[z]} &= e^{-2i\theta} \left((1 + z\bar{z})\partial_z + \frac{i}{2}\bar{z}\partial_\theta \right), \\ \nabla_{(-)}^{[z]} &= e^{2i\theta} \left((1 + z\bar{z})\partial_{\bar{z}} - \frac{i}{2}z\partial_\theta \right).\end{aligned}$$

around the north pole of sphere.

Description include south pole is achieved by

$$z = 1/w, \quad \text{and} \quad \theta' = \theta + \arg(z) + \frac{\pi}{2}, \quad \text{Then}$$

$$\begin{aligned}\nabla_{(+)}^{[w]} &= e^{-2i\theta'} \left((1 + w\bar{w})\partial_w + \frac{i}{2}\bar{w}\partial_{\theta'} \right), \\ \nabla_{(-)}^{[w]} &= e^{2i\theta'} \left((1 + w\bar{w})\partial_{\bar{w}} - \frac{i}{2}w\partial_{\theta'} \right),\end{aligned}$$

Generally , matrix may be expanded as like

$$A_{(a)} = a_{(a)}^{(0)} + a_{(a)}^{(1)(b)} i \nabla_{(b)} + b_{(a)}^{(1)bc} \mathcal{O}_{bc} + a_{(a)}^{(2)(b)(c)} i^2 \nabla_{(b)} \nabla_{(c)} + \dots$$

local Lorentz generator

gauge field *vielbein* *spin connection* *higher spin fields*

There are Infinite number of fields with various spin.

Applications of the HKK : “ $A \sim \nabla$ ”

Driving

1. Classical gravity equation

2. Quantum gravity equation

from the YM- matrix model : $S = -\text{Tr} [A,A][A,A]/4$

Classical Gravity Equations from the Matrix Model

Matrix equation of motion:

$$[A^{(b)}, [A_{(b)}, A_{(a)}]] = 0$$

***provides equations for infinitely many fields
with various spins .***

→ truncate only $A \rightarrow \nabla_a = e_a^b \partial_b + \omega_a^{bc} \mathcal{O}_{bc}$

vielbein: e_μ^a , spin connection: ω_a^{bc}

Gravity Equations

(Furuta, Hanada, Kawai, Kimura [06])

$$\begin{aligned} R_{ab} + \nabla^c T_{b,ca} + T^{p,q}{}_a T_{b,pq} &= 0, \\ \nabla^a R_{abcd} + T^{p,q}{}_b R_{pqcd} &= 0, \end{aligned}$$

torsion: $-T^c{}_{,ab} = e_\mu{}^c (\partial_a e^\mu{}_b - \partial_b e^\mu{}_a) + \omega_{ab}{}^c - \omega_{ba}{}^c,$

curvature: $R_{ab}{}^{cd} = \partial_a \omega_b{}^{cd} - \partial_b \omega_a{}^{cd} + \omega_a{}^c{}_e \omega_b{}^d{}_e - \omega_b{}^c{}_e \omega_a{}^d{}_e$

Including torsion degrees of freedom

Classical solutions (H. Isono, D.T [09])

Time-dependent solutions with torsion:

- *FRW type metric ansatz*
- *Expansion with acceleration*
 - *application to dark energy?*

Spherical symmetric solution with torsion

- *Black hole solutions*
- *All founded solutions are*
asymptotically non-flat

Quantum Gravity Equation from Matrix Model

Schrodinger equation of matrix quantum mechanics:

$$i\frac{\partial}{\partial t}\Psi(A, t) = H(A, \frac{\partial}{\partial A})\Psi(A, t)$$

After HKK: "A ~ ∇",

$$i\frac{\partial}{\partial t}\Psi(\nabla, t) = H(\nabla, \frac{\partial}{\partial \nabla})\Psi(\nabla, t)$$

We get a quantum gravity equation.

Explicit form

$$i\frac{\partial}{\partial t}\Psi = \int d^d x \sqrt{h_d} \operatorname{tr}_{\hat{g}} \left\{ \frac{1}{\sqrt{h_d}} \frac{1}{2} \left(\frac{\delta}{\delta e_i^K} \partial_K \right) \sqrt{h_d} \left(\frac{\delta}{\delta e_i^L} \partial_L \right) + \frac{1}{2} \left(\frac{\delta}{\delta \omega_i^{jk}} \mathcal{O}_{jk} \right) \left(\frac{\delta}{\delta \omega_i^{lm}} \mathcal{O}_{lm} \right) \right. \\ \left. - \frac{1}{\sqrt{h_d}} \frac{1}{4} (T_{ij}{}^K \partial_K) \sqrt{h_d} (T_{ij}{}^L \partial_L) - \frac{1}{4} (R_{ij}{}^{kl} \mathcal{O}_{kl}) (R_{ij}{}^{mn} \mathcal{O}_{mn}) \right\} \Psi.$$

(T. Matsuo, W-Y. Wen, S. Zeze, D.T [08])

Covariant only in spatial directions, no boost invariance

metric: $-dt^2 + h_d(t, x)_{IJ} dx^I dx^J$

Divergent quantities appear. Some regularization is required.

Minisuperspace Model

$$ds^2 = -dt^2 + \frac{1}{y(t)^2} dx^2.$$

Quantum mechanics of $y(t)$

Schrodinger equation was explicitly solved.

A wave packet solution was constructed

$$\Psi(t, y) = \sqrt{\frac{(2\beta)^{\frac{5}{3}}}{\Gamma(5/3)} \frac{y^2}{3^{\frac{7}{6}} (\beta + iAt)^{\frac{5}{3}}}} \exp \left[-\frac{y^3}{9(\beta + iAt)} \right].$$

This wave function describes a shrinking universe in $|t| \rightarrow \infty$

$$\langle g_{xx} \rangle = \langle y^{-2} \rangle = \frac{1}{\Gamma(\frac{5}{3})} \left(\frac{2}{9} \right)^{\frac{2}{3}} \left[\frac{\beta}{\beta^2 + A^2 t^2} \right]^{\frac{2}{3}}.$$

Number of questions have been left

- *Full covariance? ----- IKKT model may give a covariant equation. But to do this, we have to derive a sort of the Schrodinger equation from the model without time.*
- *What is a good regularization? ----- by introducing finite size matrix effect?, imposing some gauge condition?*
- *How go proceed further beyond the mini-superspace model ?*

Summary and Outlook

Various spacetime dynamics are studied by IKKT matrix model

Finite size matrix : *Non-commutative space*

Well defined, Gravity is not-manifest.

Stability analysis through the effective action

Infinite size matrix: *HKK prescription* $A \sim \nabla$

More formal, Gravity is manifest.

We got quantum and classical equations of a torsion gravity

*We need more understanding on large N
(matrix size $\rightarrow \infty$) limit.....*

*Matrix Renormalization Group (matrix RG) may be a
useful approach.*

Integrating out a part of matrix \rightarrow flow equation

*Links to the critical phenomena, Wilsonian RG are
also expected.*

(END)