

Dynamics of nonlinear wave-packet: Soliton, Vortex, Optical Patterns

YuanYao Lin

Institute of Photonics Technologies, National Tsing Hua University, Hsinchu City,
300 Taiwan



Outline

1

Introduction

- Linear/Nonlinear Waves
- About Solitons

2

Current Solitons Research

- Nonlocal Solitons
- Nonlinear Surface Plasmonics
- Soliton and Vortex in elliptical geometry
- Optical Pattern Formation: From MI to TI
- Ultrashort Solitons

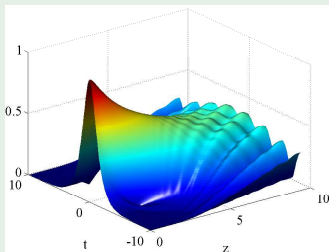
3

Conclusion



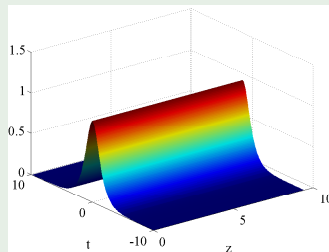
Linear/Nonlinear Wave-Packet

Linear Wave propagation



- Dispersive (time, frequency)
- Diffractive (space, momentum)

Non-Linear Wave propagation

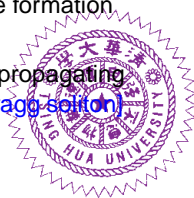


- Temporal localization
- Spatial localization
- Optical Bullet



Fundamentals, instability, transverse instability, oscillatory instability

- Solitons are nonlinear self-guided waves that propagates in a *shape (envelope) preserving* manner.
- Soliton can be bright, dark or gray depending on the form of nonlinearity.
- Modulation instability, transverse modulation instability are important menifesitation to strong nonlinear systems.
- *Modulation instability*: what soliton formation results from.
- When dimension increases, *transverse instability* govern the formation of 2D solitary wave structure. [\[Quasi-1D soliton stripe\]](#)
- When there are multiple components, for example, counter propagating fields, *oscillatory instability* from nonlinear interferences. [\[Bragg soliton\]](#)



Story: First observation of (water surface) soliton

Scottish engineer, **John Scott Russell** (1808-1882), "Report on Waves": **The wave of translation**, (*Report of the fourteenth meeting of the British Association for the Advancement of Science, York, September 1844 (London 1845), pp 311-390, Plates XLVII-LVIII.*)

First Observation- Waves of Translation



John Scott Russell
(1808-1882)

Such, in the month of August 1834, was my first chance interview with that singular and beautiful phenomenon which I have called the Wave of Translation.



Universal signature of solitons

A Universal phenomenon of self-trapped wave packets

- EM waves in nonlinear optical materials;
- shallow- and deep-water waves;
- charge-density waves in plasmas;
- matter waves in Bose-Einstein condensates;
- excitations on DNA chains;
- domain walls in supergravity, and "branes" at the end of open strings in superstring theory; to name only a few.

See: M. Segev, *Optics & Photonics News*, pp. 27 (Feb. 2002).



Current Studies on Solitons

In the study of solitons, what is the question in concern currently?

- What is the Applications?
soliton based communication system, optical switching, and the like. [Boardman A D and Sukhorukov A P (ed) 2001 *Soliton-Driven Photonics* (Dordrecht: KluwerAcademic)]
- What is the material?
nonlocal response: thermal-optical (lead glass), photorefractive(SBN), liquid crystal, resonance absorbing media. [A. W. Synder and D. J. Mitchell, "Accessible Solitons," *Science* **276**, 1538-1541 (1997).]
- What is the boundary effect?
gap soliton in (subwavelength)lattice, **surface waves**. [de Sterke C M and Sipe J E 1994 Gap solitons Progress in Optics vol XXXIII ed E Wolf (Amsterdam: North-Holland) pp 20360; Mingaleev S and Kivshar Yu S 2001 Phys. Rev. Lett. 86 5474; Slusher R and Eggleton B (ed) 2003 *Nonlinear Photonic Crystals* (Berlin: Springer); A. V. Zayats, I. I. Smolyaninov, and A. A. Maradudin, *Phys. Rep.* **408**, 131 (2005).]
- Dissipative or Conservative?
cavity soliton, laser soliton(autosoliton) [N. Akhmediev and A. Ankiewicz (ed), *Dissipative Solitons: From Optics to Biology and Medicine*, Lect. Notes Phys. 751 (Springer, Berlin Heidelberg (2008).]
- What is the time scale?
non-instantaneous, ultrashort pulse [P. Kinsler and G.H.C. New, "Few-cycle soliton propagation," *Phys. Rev. A* **69**, 013805 (2004).]



Nonlocal Response

Nonlinear response related to the neighboring materials

$$i \frac{\partial U}{\partial z} + \frac{1}{2} \nabla_{\perp}^2 U + nU = 0, \quad n = \int_{-\infty}^{\infty} R(\eta - \eta') |U(\eta')|^2 d\eta'$$

Diffusive response

$$n - d \nabla_{\perp}^2 n = |U|^2$$

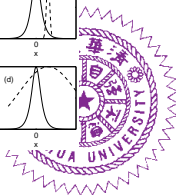
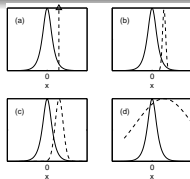
$$\Downarrow$$

$$R(\eta) = \frac{1}{\sqrt{2d}} e^{-|\eta|/\sqrt{d}}$$

photorefractive material
thermal-optical material
nematic liquid crystal
plasma

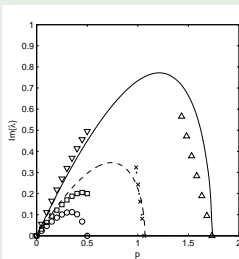
Dipole-dipole force in dipolar BEC

$$R(\vec{\eta}, \vec{\eta}') = \frac{1 - 3\vec{p} \cdot \vec{e}_{\vec{\eta} - \vec{\eta}'}}{|\vec{\eta} - \vec{\eta}'|^3}$$

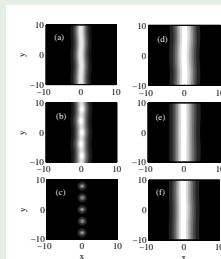


Transverse instability suppression in nonlocal nonlinear medium

TI spectrum



nonlinear wave evolution



Nonlocal nonlinear response can suppress soliton transverse instabilities, shown by numerical calculation and asymptotic expansions.

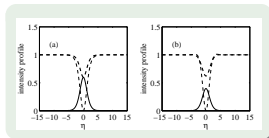
See: Y.Y. Lin et al., *J. Opt. Soc. Am. B* **25**, 576-581 (2008)

Dark-bright soliton pair in nonlocal nonlinear media

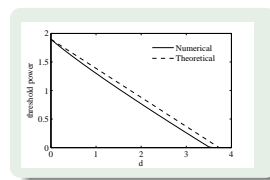
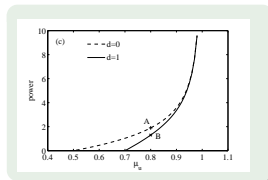
typical nonlocal response \iff formation power increase !

$$i\partial_{\eta}\psi - \frac{1}{2}\partial_{\xi}\xi + n\psi = 0, \quad \psi = U(\xi; \eta), \quad V(\xi; \eta)$$

$$n = \int_{-\infty}^{\infty} R(\xi' - \xi) (|U|^2 + |V|^2), \quad R(\xi) = \frac{1}{2\sqrt{d}} e^{-|\xi|/\sqrt{d}}$$



(a): $d=0$ and (b): $d=1$.



when bright pulse is coupled to dark pulse in nonlocal response \iff
formation power decrease !

See: Y.Y. Lin et al., *Opt. Express* **15**, 8781-8786 (2007)

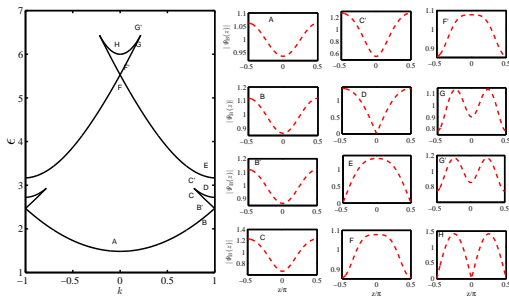
Nonlinear bloch wave and swallow-tailed band structure

Nonlinear band diagram with dipolar nonlocal response

$$\mu \Psi(\vec{r}) = \left[-\frac{\hbar^2}{2m} \nabla^2 + V_{\text{trap}} + Ng_s |\Psi(\vec{r})|^2 + Ng_d \int d^3 r' V_d(\vec{r}, \vec{r}') |\Psi(\vec{r}')|^2 \right] \Psi(\vec{r}), \quad \frac{1 - 3\vec{p} \cdot \vec{e}_{r-r'}}{|\vec{r} - \vec{r}'|^3}$$

- Cigar shaped dipolar BEC in optical lattice.
- Effective 1D nonlocal model.
- Sufficient Bloch wave expansion up to 3-modes.

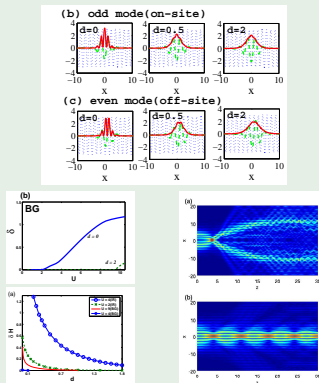
Band structure manipulation by nonlocality control, which is implemented with dipole orientation angle.



See: Y.Y. Lin et al., Phys. Rev. A **78**, 023629 (2008).

Gap solitons and Bragg Solitons

nonlocal gap soliton



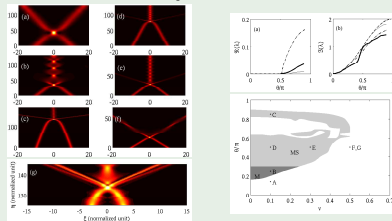
See: Y.Y. Lin et al., J. Opt. A: Pure Appl. Opt. **10** 044017 (2008).

nonlocal bragg solitons

$$i \frac{\partial U}{\partial \eta} + i \frac{\partial U}{\partial \xi} - \frac{1}{2} \left[n_0 + \frac{|V|^2}{1+Dg^2} \right] U + V = 0$$

$$i \frac{\partial V}{\partial \eta} - i \frac{\partial V}{\partial \xi} - \frac{1}{2} \left[n_0 + \frac{|U|^2}{1+Dg^2} \right] V + U = 0$$

$$n_0 - D \frac{\partial^2 n_0}{\partial \xi^2} = |U|^2 + |V|^2$$



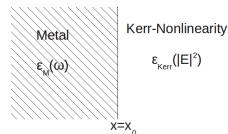
See: Y.Y. Lin et al., Phys. Rev. A, **80**, 013838 (2009).

Nonlinear Surface Plasmonics

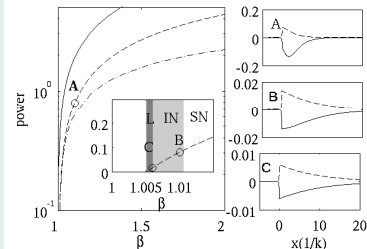
Assume the divergence of displacement vanishes,

$$\left[\nabla^2 + k^2 n^2 \right] E = -\nabla \left[E \cdot \nabla \ln n^2 \right]$$

E : vector electric field, k : wave vector in vacuum, n : scalar refractive index ($n = \sqrt{\vec{E} \cdot \vec{E}}$).

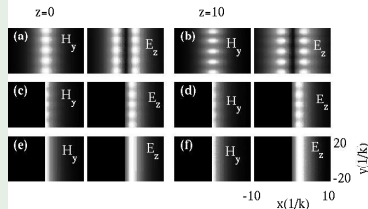


Nonlinear Surface mode



See: Y.Y. Lin et al., Opt. Lett. **34**, 2982-2984 (2009)

Stable evolution of nonlinear surface plasmons



Singular waves: Linear/Nonlinear Vortex

- A singular optical waves with "topological charges", i.e. phase change in the azimuthal direction;
- Possessing Phase singularities;
- Carrying angular momentum;
- Linear or Nonlinear.

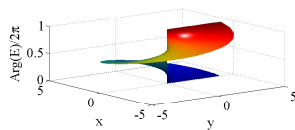
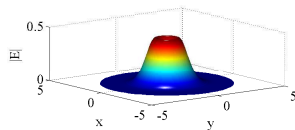
Paraxial *linear* vortex beam

$$E^{(m)}(\rho, \phi, z) = E_s \frac{\rho_s}{w_s} \left(\frac{\rho}{w_s} \right)^{|m|} \exp \left(-\frac{\rho^2}{w_s^2} \right) \exp [i\Phi_s(\rho, \phi, z)]$$

$$\Phi_s(\rho, \phi, z) = -(|m| + 1) \tan^{-1} \left(\frac{2z}{k\rho_s^2} \right) + \frac{k\rho^2}{2R_s(z)} + m\phi + kz,$$

$$w_s = \sqrt{\rho_s^2 + \left(\frac{2z}{k\rho_s} \right)^2}, \quad R_s(z) = z + \frac{k^2 \rho_s^4}{4z}.$$

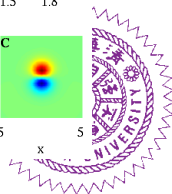
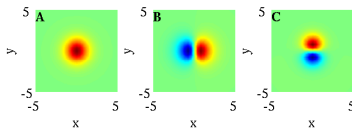
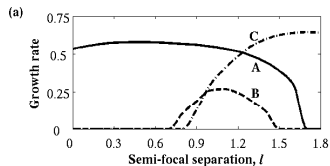
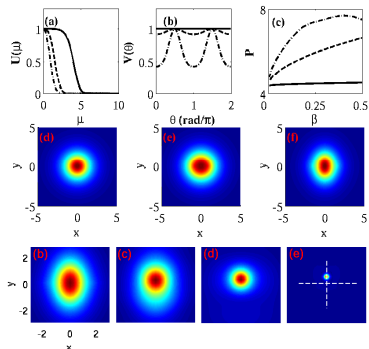
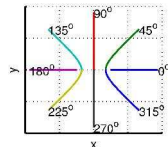
See: L. Allen et al., *Optical Angular Momentum*, Institute of physics Publishing, (2003)



Elliptical Soliton

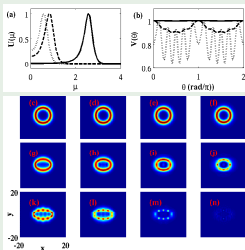
Conformal Mapping

Apply to nonlinear Schrödinger equation, we let $x = l \cdot \cosh \mu \cos \theta$ and $y = l \cdot \sinh \mu \sin \theta$, leading to $\nabla_{\perp}^2 = \frac{1}{l^2 (\cosh^2 \mu - \cos^2 \theta)} \left[\frac{\partial^2}{\partial \mu^2} + \frac{\partial^2}{\partial \theta^2} \right]$ where l is the semi-focal separation, $\mu > 0$ and $\theta \in [0, 2\pi]$.

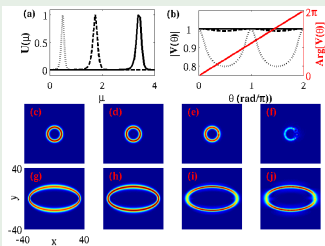


Elliptical Vortex

Ring Soliton (0π)



Vortex (2π)



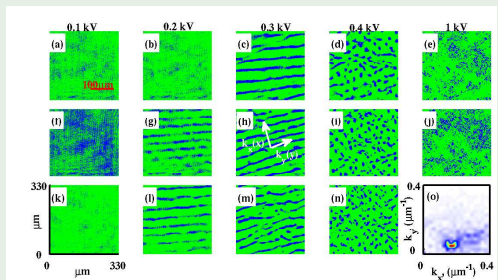
- Elliptical soliton rings evolve to clusters in a specific geometric arrangement ascribable to elliptical symmetry and the corresponding transverse instabilities.
- Elliptical vortex breaks at the major axis, turning into crescents and end up with clusters.

See: YuanYao Lin et al., Opt. Lett. **33**, 1377 (2008).

Optical Pattern in Photorefractive materials

The optical intensity patterns of a coherent beam at the output plane through a nonlinear crystal at different bias voltages and different input intensities

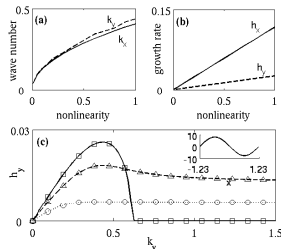
Experimental Result



(a-e): 65 $\frac{mW}{cm^2}$; (f-j): 96 $\frac{mW}{cm^2}$; (k-n): 96 $\frac{mW}{cm^2}$; (o): Fourier spectrum of (h).

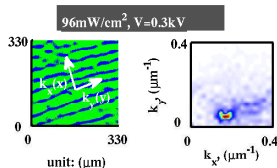
See: C.-C. Jeng, Y.Y. Lin, R.-C. Hong, and R.-K. Lee, Phys. Rev. Lett. **102** 153905 (2009)

MI/TI analysis



- Characteristic spatial vectors for MI and TI are close;
- MI growth rate \gg TI growth rate;
- numerics and analytics match well;
- spatial vector with a maximal growth rate is independent of $\Omega\tau$

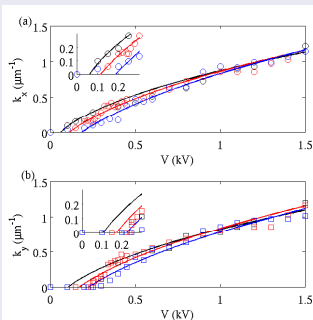
Optical Pattern in Photorefractive materials



Spatial vectors: $k_x(a)$ and $k_y(b)$

Spectra of transverse patterns obtained from experimental data for (a) k_x and (b) k_y . Circles and squares markers accompanied by the fitted curves in black, red and blue refer to optical intensities $I = 190, 96$, and 65 mW/cm², respectively.

Experimental, Theoretical and TI threshold



Supporting Facts

- second threshold voltage:
 $V_{th}^y > V_{th}^x \Leftrightarrow h_y < h_x$
- intensity dependent threshold voltage:
 $I \nearrow \Rightarrow V_{th}^y, V_{th}^x \searrow \Leftrightarrow I \nearrow \Rightarrow \tau \searrow$

I	65mW/cm ²	96mW/cm ²	190mW/cm ²
V_{th}^x	0.19kV	0.117kV	0.067kV
V_{th}^y	0.23kV	0.17kV	0.108kV

Ultrashort Solitons: Corrections non-instantaneous(nonlocal) Polarization

For solitons, we study wave packet of special envelope character within the scope of slowly varying envelope approximation (SVEA).

Under this scope inverse scattering transformation (IST), a well developed method to analyze the interaction between solitons, can be successfully applied of system integrability

When SVEA fails, IST thus fails. It requires new techniques:

- FDTD [A. Taflov, *Computational Electrodynamics: The Finite Difference Time-Domain Method*, Boston: Artech House, 1995];
- slowly evolving wave approximation(SEWA) [M.A. Porras, Phys. Rev. A, **60**,5059 (1999)];
- generalized few-cycle envelope approximation (GFEA)[P. Kinsler, Phys. Rev. A **81**, 013819 (2010)].

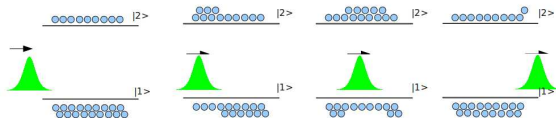
See: [arXiv:physics/0212014v4](https://arxiv.org/abs/physics/0212014v4) for details



Self-induced-transparency(SIT) solitons in Two-level-atom(TLA)

- A coherent optical pulse propagate in a resonant absorbing medium without any experiences of loss and distortion when exceed certain power.
- In particular, optical solitons in such a resonance are referred to as self-induced-transparency **SIT** soliton

Ref: S.L. McCall and E.L. Hahn, "Self-Induced Transparency," Phys. Rev. **183**, 457-489 (1969).



E : Electric field envelope of EM field;
 $N \equiv N_{|2\rangle} - N_{|1\rangle}$: population difference;
 $P \propto \rho_{12}$: transition dipole moment.

$$k_j = \frac{\partial j}{\partial \omega}, j = 1, 2, 3 \dots$$

Maxwell-Bloch Equation

$$\frac{\partial}{\partial t} P = \frac{1}{2} \frac{\mu^2}{\hbar} N E \quad ; \quad \frac{\partial}{\partial t} N = -\frac{1}{\hbar} (P E^* + P^* E);$$

$$\left[\frac{\partial}{\partial t} + c \frac{\partial}{\partial z} \right] E = \frac{\omega_0}{\epsilon_0} P + \frac{\omega_0}{\epsilon_0} \left(2 - \frac{\omega_0 k_1}{k_0} \right) \frac{i}{\omega_0} \frac{\partial}{\partial t} P;$$

Ultrashort SIT solitons

Taking all the arguments and constraints, we can derive a single **cubic-quintic** nonlinear ordinary differential equation(ODE) from the corrected **MB** equation that governs the solitary wave solution.

cubic-quintic equation for $q = |E|$

$$\frac{\partial^2 q}{\partial \xi^2} - \left(\frac{\Gamma n_c}{4\alpha} - \beta^2 \right) q + \left(\frac{1}{2} - \frac{\bar{\sigma}\beta}{2} \right) q^3 + \frac{3\bar{\sigma}^2}{64} q^5 = 0$$

in which $\beta \equiv -\frac{\Delta}{2\alpha} + \frac{\Gamma \bar{\sigma} n_c}{8\alpha}$

temporal phase $\phi(\xi) = \angle E$

$$\frac{\partial \phi}{\partial \xi} = -\frac{3}{8} \bar{\sigma} q^2 + \beta$$

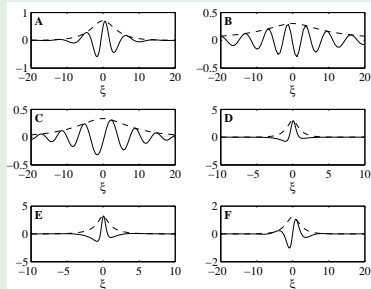
atomic polarizations n , $u = \Re\{P\}$, and $v = \Im\{P\}$

$$n = n_c - \frac{\alpha}{\Gamma} q^2$$

$$u = \frac{2\alpha}{\Gamma} \frac{\partial q}{\partial \xi}$$

$$v = \frac{2\alpha}{\Gamma} \left(\frac{\bar{\sigma}}{8} q^3 - \beta q \right)$$

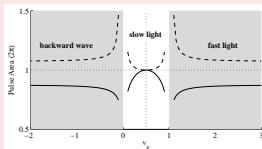
Soliton solutions to $\Re\{E\}$ (dashed) and $|E|$ (solid) at $\Delta = 0$



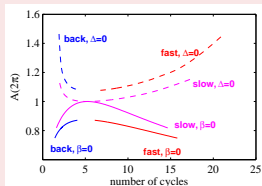
A: $n_c = 1, v_g = -1$; **B:** $n_c = 1, v_g = -0.1$;
C: $n_c = -1, v_g = 0.1$; **D:** $n_c = -1, v_g = 0.9$;
E: $n_c = 1, v_g = 1.1$; **F:** $n_c = 1, v_g = 3$.

Stable Backward wave, slow light, superluminal and Area theory breakdown

Area-velocity



Area-Pulse duration



General Soliton family and area theory

- Soliton solution exists only when $n_c \alpha > 0$.
- **slow light soliton**: $n_c < 0$, which means atoms are more populated in the ground state $|1\rangle$, the group velocity of the soliton $v_g < 1$ (normalized to the velocity of light, c).
- **fast light/superluminal soliton**: $n_c > 0$, which means more atoms are excited to $|2\rangle$, the group velocity of the soliton can be $v_g > 1$ (group velocity exceeds the velocity of light, c) Phys. Rev. A **75**, 043806(2007).
- **backward wave soliton**: $n_c > 0$, which means more atoms are excited to $|2\rangle$, the group velocity of the soliton can be $v_g < 0$ (propagating in the reverse direction with respect to phase velocity).
- stability confirmed by linear stability analysis
- **dark solitons** are not possible due to violation of conserved quantity for Bloch vectors.
- **Area theory** breakdown is illustrated in **real ultrashort SIT solitons**.

conclusion

- Briefly, concepts to solitons is introduced.
- Swiftly, we browse some of the current study associated with solitons: nonlocal effect, special geometry, conserved/dissipative, non-instantaneous effect, and ultrashort soliton.
- Hopefully, interesting physics can be explored from these studies.



Acknowledgement

Colaborators

Prof. Yuri S. Kivshar, Nonlinear Physics Center, Research School of Physical Sciences and Engineering, The Australian National University, Canberra, ACT 0200, Australia

Prof. Boris A. Malomed, Department of Physical Electronics, School of Electrical Engineering, Faculty of Engineering, Tel Aviv University, Tel Aviv 69978, Israel.

Dr. T.-D. Lee, Industrial Technology Research Institute, Hsinchu, 310 Taiwan

Dr. J.-S. Pan, TrueLight Corporation, Hsinchu, 300 Taiwan

Prof. C.-C. Jeng, Department of Physics, National Chung-Hsing University, Taichung 402, Taiwan

Funding Agents

The author are indebted to NSC for the funding support. Contract NO.

98-2112-M-007-012

Thank you

